

ON UNIQUENESS OF SOLUTIONS FOR NONLINEAR CREEP PROBLEMS IN SYMMETRIC PRESSURE VESSELS†

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Abstract—Uniqueness is established for positive solutions of a nonlinear integral equation which governs the effective stress in internally loaded spherical and incompressible cylindrical pressure vessels subject to primary or secondary transient creep.

INTRODUCTION

It was shown in [1, 2] that the boundary value problems for transient quasistatic creep in hollow spherical and infinite incompressible cylindrical pressure vessels subject to a radial nondecreasing internal pressure can be reduced to an integral equation of the form

$$\sigma(r, t) = \frac{\beta P(t)}{r^j} + \frac{\mu \beta}{r^j} \int_a^b \left[\int_0^t \sigma^n(\xi, \tau) d\tau \right]^{1/(m+1)} \frac{d\xi}{\xi} - \mu \left[\int_0^t \sigma^n(r, \tau) d\tau \right]^{1/(m+1)}. \quad (1)$$

Here, the unknown function σ stands for the effective stress, and $P(t)$ is proportional to the prescribed internal pressure. The independent variables r and t denote, respectively, radial distance from the center of the vessel and time, a and b are the internal and external radii of the vessel, and

$$\beta^{-1} \equiv \int_a^b \frac{d\xi}{\xi^{j+1}} \quad (2)$$

where $j = 2$ for cylinders, $j = 3$ for spheres.

The quantities n and m are creep constants arising from the well-known [3] primary creep law

$$\epsilon_{ij}^{(c)} = \frac{3K}{2} \int_0^t \frac{\sigma_e^{n-1}}{[\epsilon_e^{(c)}]^m} s_{ij} d\tau \quad (3)$$

relating the creep strains $\epsilon_{ij}^{(c)}$ to the effective creep strain‡ $\epsilon_e^{(c)}$, the effective stress σ_e , and the deviatoric stress components s_{ij} . Notice that, for $m = 0$, (3) reduces to the law for secondary creep, and the structure of (1) is greatly simplified. The constant μ depends on both elastic and creep constants and is proportional to $K^{1/(m+1)}$. Thus, in the absence of creep ($K = 0$), $\mu = 0$ and eqn (1) collapses to the solution of the elastic problem.

For the present analysis and in order to be able to use the results of [1] and [2] on boundedness of solutions of (1), we must assume that

$$0 \leq m \leq n - 1, \quad \mu > 0 \quad (4)$$

$$0 < P(0), \quad 0 \leq \dot{P}(t) \quad (t > 0) \quad (5)$$

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‡See [1] for detailed definitions of $\epsilon_e^{(c)}$, etc.

and that \dot{P} is integrable on $[0, \infty]$. Also, for $m > 0$, it is assumed that

$$\sigma(r, t) > 0 \quad (a \leq r \leq b, t > 0). \quad (6)$$

In the case $m = 0$, (6) can be shown to be a consequence of (4) and (5).

With the above assumptions, it was shown in [2] for secondary creep and in [1] for primary creep that any positive solution σ of (1) has the following bounds:

$$\beta \frac{P(t)}{b^j} \leq \sigma(r, t) \leq \beta \frac{P(t)}{a^j} \quad (a \leq r \leq b, t \geq 0). \quad (7)$$

It is the purpose of this note to show that there is only one such solution. The main idea of the proof is a uniqueness argument based on the Gronwall-type inequality of Lemma 2 below. This line of attack was suggested by the work of Lewis Wheeler in [4] on the uniqueness of solutions to the displacement problem of nonlinear dynamic elasticity.

UNIQUENESS THEOREM

We shall need the following two elementary lemmas.

LEMMA 1. For x_1, x_2 , and m positive,

$$|x_1^{1/(m+1)} - x_2^{1/(m+1)}| \leq \frac{1}{m+1} (\min\{x_1, x_2\})^{-m/(m+1)} |x_1 - x_2|. \quad (8)$$

A simple proof of this inequality can be obtained by using a finite Taylor expansion of the function $f(x) \equiv x^{1/(m+1)}$ about the smaller of x_1 and x_2 and the fact that $f'' < 0$ on $(0, \infty)$.

Lemma 2. (Gronwall). Let $u(t)$ be continuous on $[0, T]$ and suppose that, for some finite constants $C > 0$ and $0 \leq \alpha < 1$,

$$0 \leq u(t) \leq Ct^{-\alpha} \int_0^t u(\tau) d\tau \quad (0 < t \leq T). \quad (9)$$

Then

$$u(t) = 0 \quad (0 \leq t \leq T). \quad (10)$$

This lemma is proved by using an integrating factor to put (9) in the form

$$\frac{d}{dt} \left\{ \left(\int_0^t u(\tau) d\tau \right) \exp \left[\frac{Ct^{1-\alpha}}{\alpha-1} \right] \right\} \leq 0.$$

Theorem. The integral eqn (1) has at most one continuous, positive solution σ on $[a, b] \times [0, \infty)$.

Proof. Suppose that σ_1 and σ_2 are both positive solutions of (1) and define

$$\bar{\sigma}(r, t) \equiv \sigma_1(r, t) - \sigma_2(r, t). \quad (11)$$

Then, it suffices to show that

$$u^2(t) \equiv \int_a^b r^{j-1} \bar{\sigma}^2(r, t) dr = 0 \quad (12)$$

on any interval $[0, T]$, $T > 0$.

It follows from (1) and (11) that

$$\begin{aligned} \bar{\sigma}(r, t) = & \frac{\mu\beta}{r^j} \int_a^b \left(\left[\int_0^t \sigma_1^n(\xi, \tau) d\tau \right]^{1/(m+1)} - \left[\int_0^t \sigma_2^n(\xi, \tau) d\tau \right]^{1/(m+1)} \right) \frac{d\xi}{\xi} \\ & - \mu \left(\left[\int_0^t \sigma_1^n(r, \tau) d\tau \right]^{1/(m+1)} - \left[\int_0^t \sigma_2^n(r, \tau) d\tau \right]^{1/(m+1)} \right). \end{aligned} \quad (13)$$

If this equation is now multiplied by r^{-1} and integrated from a to b , it will follow, using (2), that

$$\int_a^b \bar{\sigma}(r, t) \frac{dr}{r} = 0 \quad (0 \leq t < \infty). \tag{14}$$

Because of (14), if we multiply both sides of (13) by $r^{j-1} \bar{\sigma}$ and integrate from a to b , we shall obtain

$$u^2(t) = -\mu \int_a^b r^{j-1} \bar{\sigma}(r, t) \left(\left[\int_0^t \sigma_1^n(r, \tau) d\tau \right]^{1/(m+1)} - \left[\int_0^t \sigma_2^n(r, \tau) d\tau \right]^{1/(m+1)} \right) dr. \tag{15}$$

Notice that, by virtue of (5) and (7),

$$\min \left\{ \int_0^t \sigma_1^n(r, \tau) d\tau, \int_0^t \sigma_2^n(r, \tau) d\tau \right\} \geq \left[\frac{\beta P(0)}{b^j} \right]^n t. \tag{16}$$

Therefore, we can use (16) and Lemma 1, taking

$$x_i = \int_0^t \sigma_i^n(r, \tau) d\tau \quad (i = 1, 2),$$

to derive from (15) the inequality

$$u^2(t) \leq \frac{\mu}{m+1} \left(\frac{\beta P(0)}{b^j} \right)^{-mn/(m+1)} t^{-mn/(m+1)} \int_a^b r^{j-1} |\bar{\sigma}(r, t)| \int_0^t |\sigma_1^n(r, \tau) - \sigma_2^n(r, \tau)| d\tau dr. \tag{17}$$

Due to (7) and the continuity of P , it is clear that there exists a finite constant $M = M(T)$ such that

$$|\sigma_1^n(r, \tau) - \sigma_2^n(r, \tau)| \leq M |\bar{\sigma}(r, \tau)| \quad (0 \leq \tau \leq T). \tag{18}$$

In fact, M can even be chosen independent of T if P is bounded. Therefore, if we set

$$C = \frac{\mu M}{m+1} \left(\frac{\beta P(0)}{b^j} \right)^{-mn/(m+1)}, \quad \alpha = \frac{m}{m+1}, \tag{19}$$

it will follow from (17) and (18) that

$$u^2(t) \leq Ct^{-\alpha} \int_a^b \int_0^t r^{j-1} |\bar{\sigma}(r, t)| |\bar{\sigma}(r, \tau)| dr d\tau. \tag{20}$$

Therefore, by Schwarz's inequality,

$$\begin{aligned} u^2(t) &\leq Ct^{-\alpha} \left[\int_a^b r^{j-1} \bar{\sigma}^2(r, t) dr \right]^{1/2} \int_0^t \left[\int_a^b r^{j-1} \bar{\sigma}^2(r, \tau) dr \right]^{1/2} d\tau \\ &\leq Ct^{-\alpha} u(t) \int_0^t u(\tau) d\tau. \end{aligned}$$

This result, together with Lemma 2, completes the proof.

REFERENCES

1. W. S. Edelman, On bounds for primary creep in symmetric pressure vessels. *Int. J. Solids Structures* **12**, 107 (1976).
2. W. S. Edelman and R. A. Valentin, On bounds and limit theorems for secondary creep in symmetric pressure vessels. *Int. J. Non-Linear Mech.* **11**, 265 (1976).
3. F. K. G. Odqvist and J. Hult, *Kriechfestigkeit Metallischer Werkstoffe*. Springer, Berlin (1962).
4. Lewis Wheeler, *A Uniqueness Theorem for the Displacement Problem in Finite Elastodynamics*. To appear.